

First-Passage Statistics of Extreme Values

Eli Ben-Naim

Los Alamos National Laboratory

with:

Paul Krapivsky (Boston University)

Nathan Lemons (Los Alamos)

Pearson Miller (Yale, MIT)

Talk, publications available from: <http://cnls.lanl.gov/~ebn>

APS March Meeting, Baltimore MD, March 17, 2016

Plan

- I. **Motivation:** records & their first-passage statistics as a data analysis tool
- II. **Ordered records:** uncorrelated random variables
- III. **Ordered records:** correlated random variables

I. Motivation:

records & their first-passage
statistics as a data analysis tool

Extreme value statistics

New frontier in nonequilibrium statistical physics

- Brownian motion Comtet, Majumdar, Krug, Redner
- Surface growth Spohn, Halpin-Healy, Majumdar, Schehr
- Transport Mallick, Krapivsky, Derrida, Lebowitz, Speer
- Population dynamics Kamenev, Meerson, Doering, Nelson
- Climate Bunde, Havlin, Krug, Wergen, Redner
- Earthquakes Davidesn, Sornette, Newman, Turcotte, EB
- Finance Bouchaud, Stanley, Majumdar

Record & Running Record



- Record = largest variable in a series

$$X_N = \max(x_1, x_2, \dots, x_N)$$

- Running record = largest variable to date

$$X_1 \leq X_2 \leq \dots \leq X_N$$

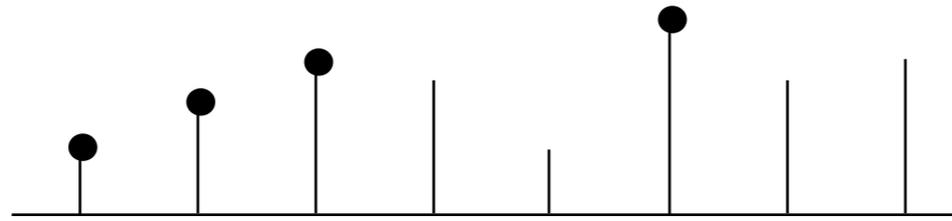
- Independent and identically distributed variables

$$\int_0^{\infty} dx \rho(x) = 1$$

Statistics of extreme values

Feller 68
Gumble 04
Ellis 05

Average number of running records



- Probability that N th variable sets a record

$$P_N = \frac{1}{N}$$

- Average number of records = harmonic number

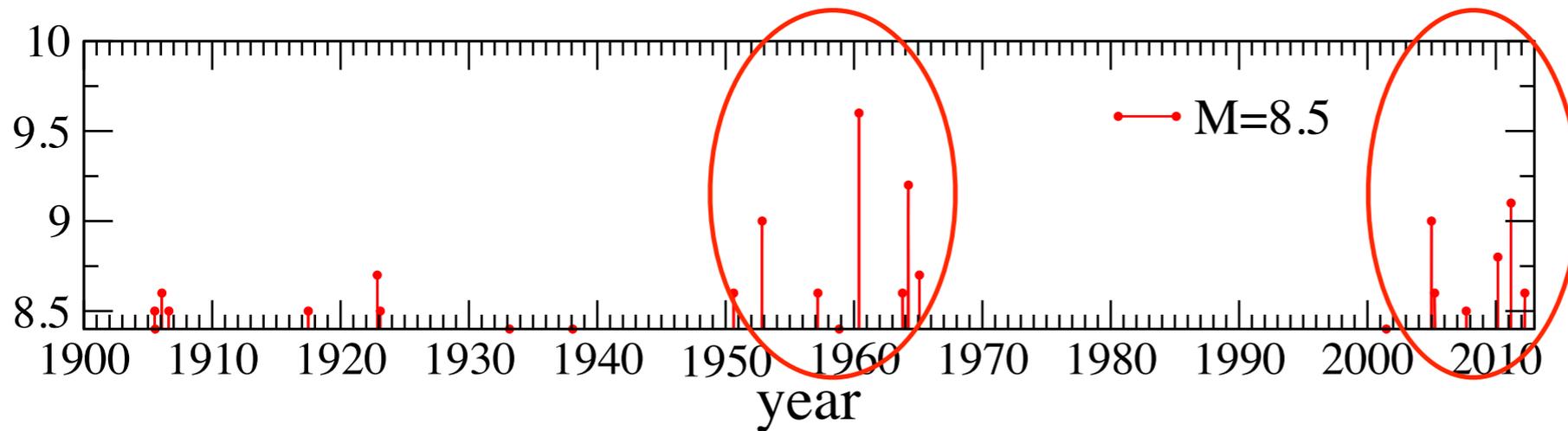
$$M_N = 1 + \frac{1}{2} + \frac{1}{3} + \cdots + \frac{1}{N}$$

- Grows logarithmically with number of variables

$$M_N \simeq \ln N + \gamma \quad \gamma = 0.577215$$

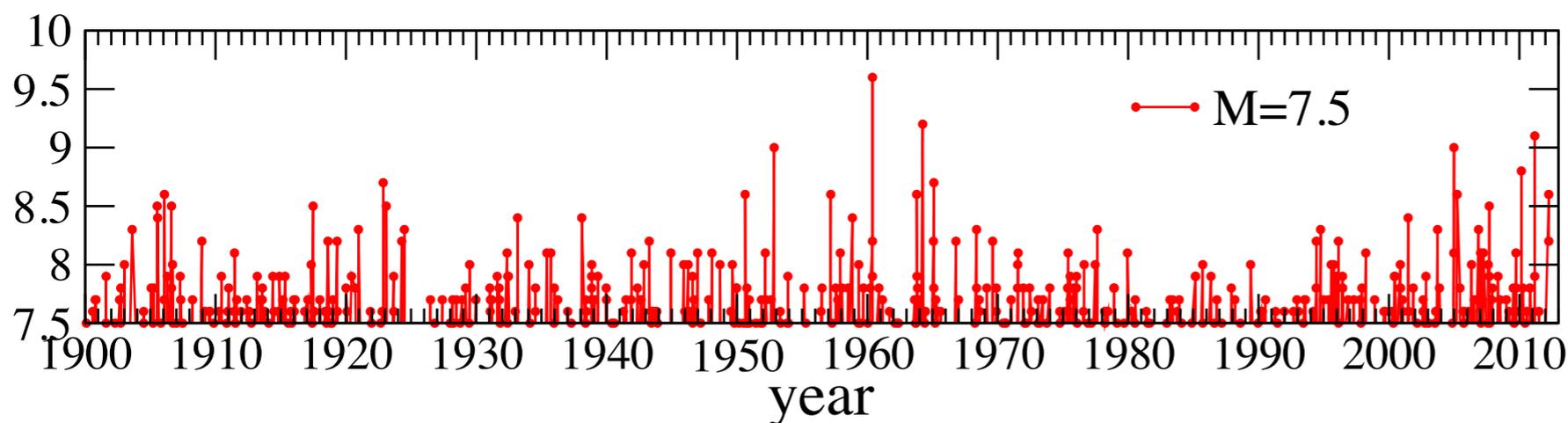
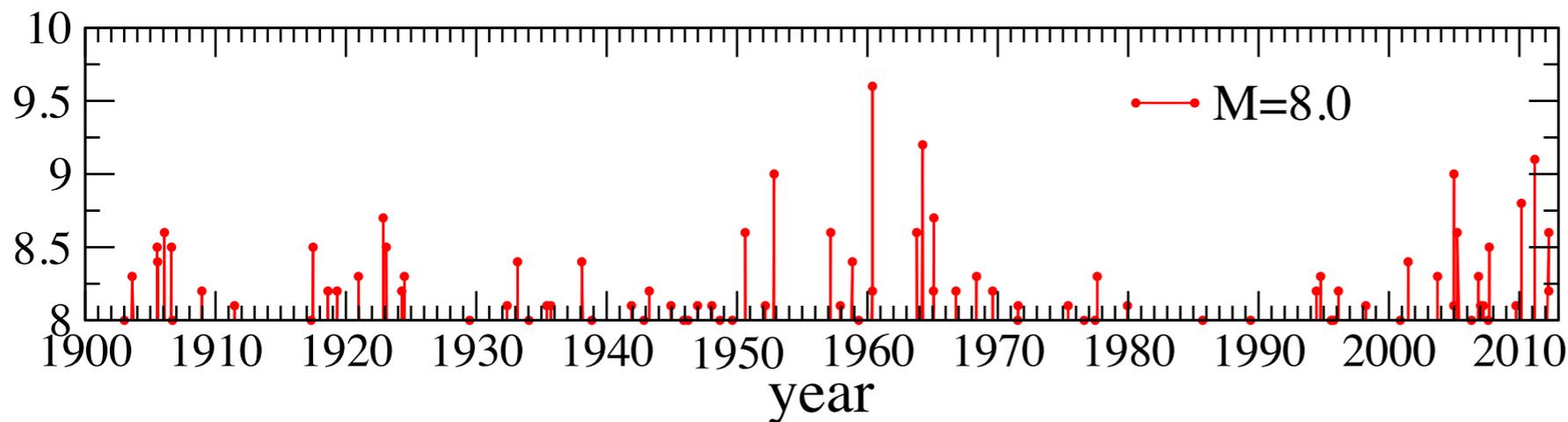
Behavior is independent of distribution function
Number of records is quite small

Clustering of massive earthquakes?



1770 $M>7$ events
1900-2013

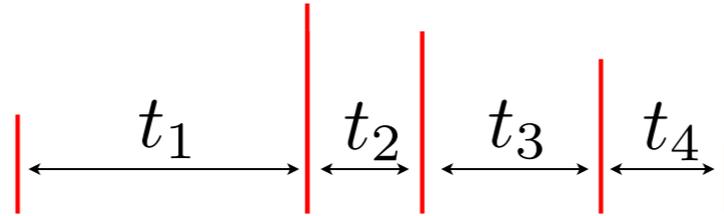
Magnitude	Annual #
9-9.9	1/20
8-8.9	1
7-7.9	15
6-6.9	134
5-5.9	1300
4-4.9	~13,000
3-3.9	~130,000
2-2.9	~1,300,000



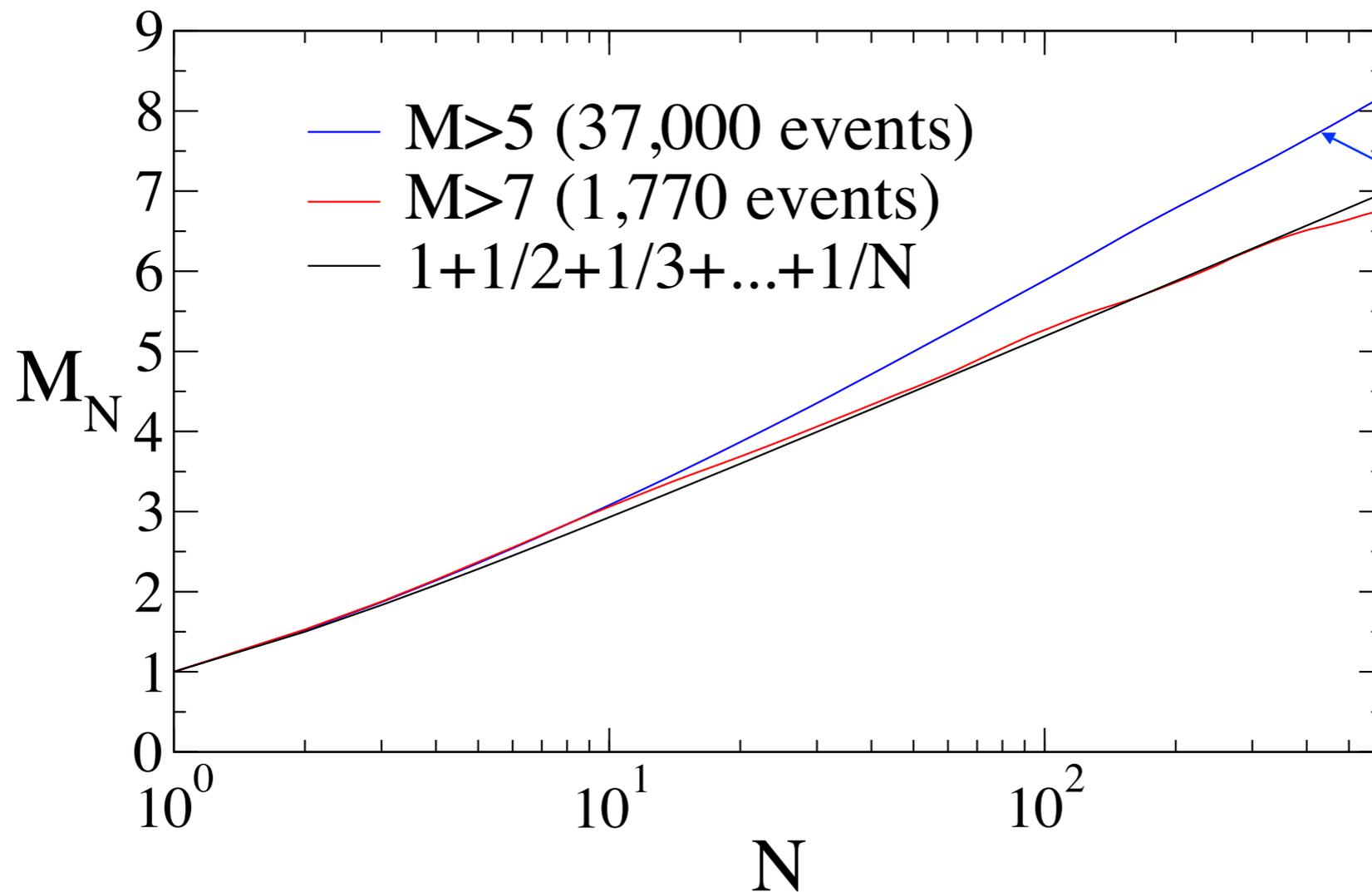
Are massive earthquakes correlated?

Earthquake Triggering
Gomberg 05
Lay 10

Records in inter-event time statistics



Count number of running records in N consecutive events



attribute
deviation to
aftershocks

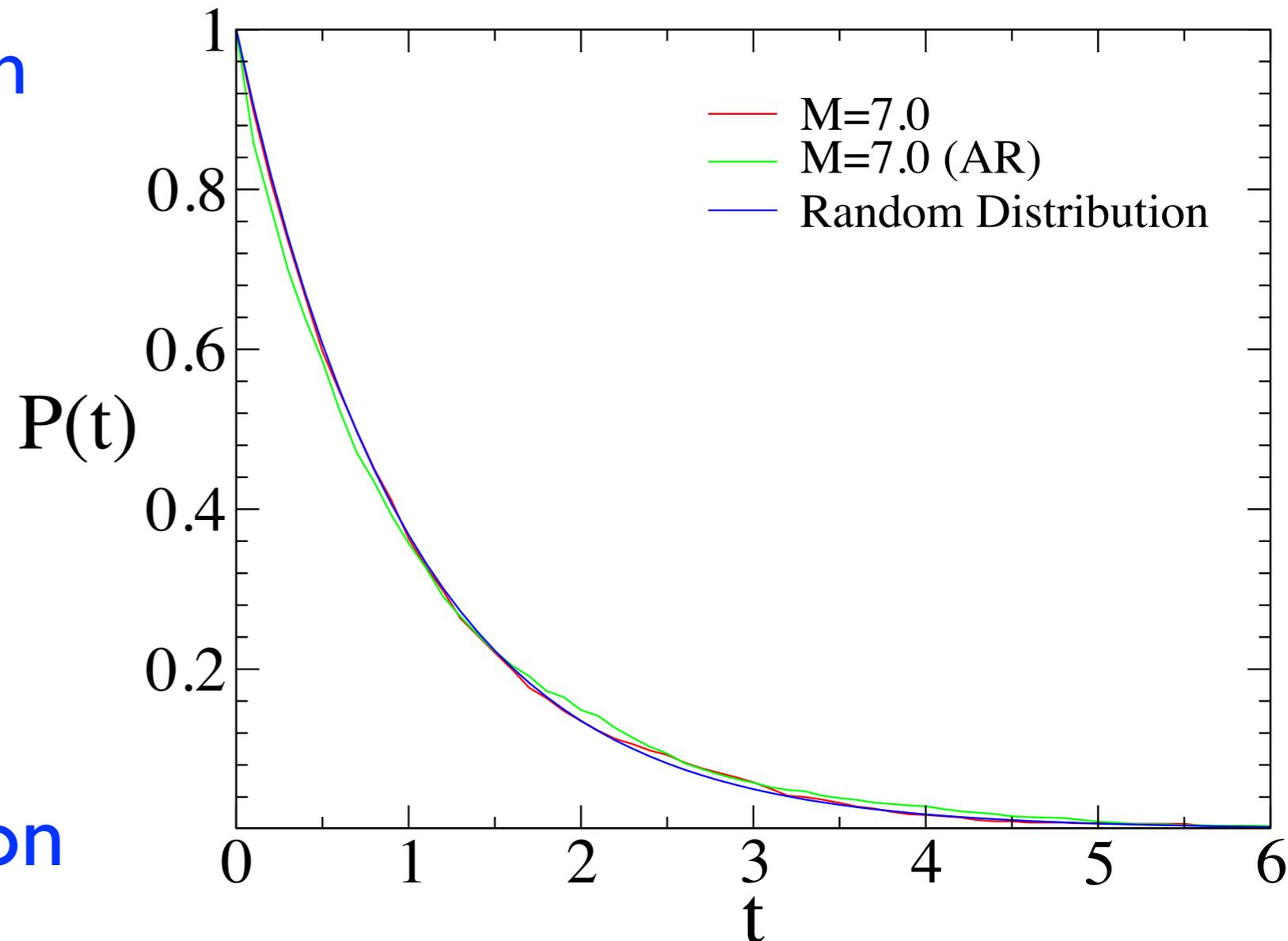
records indicate inter-event times uncorrelated

Massive earthquakes are random

Inter-Event time statistics

EB, Daub, Johnson, GRL 13

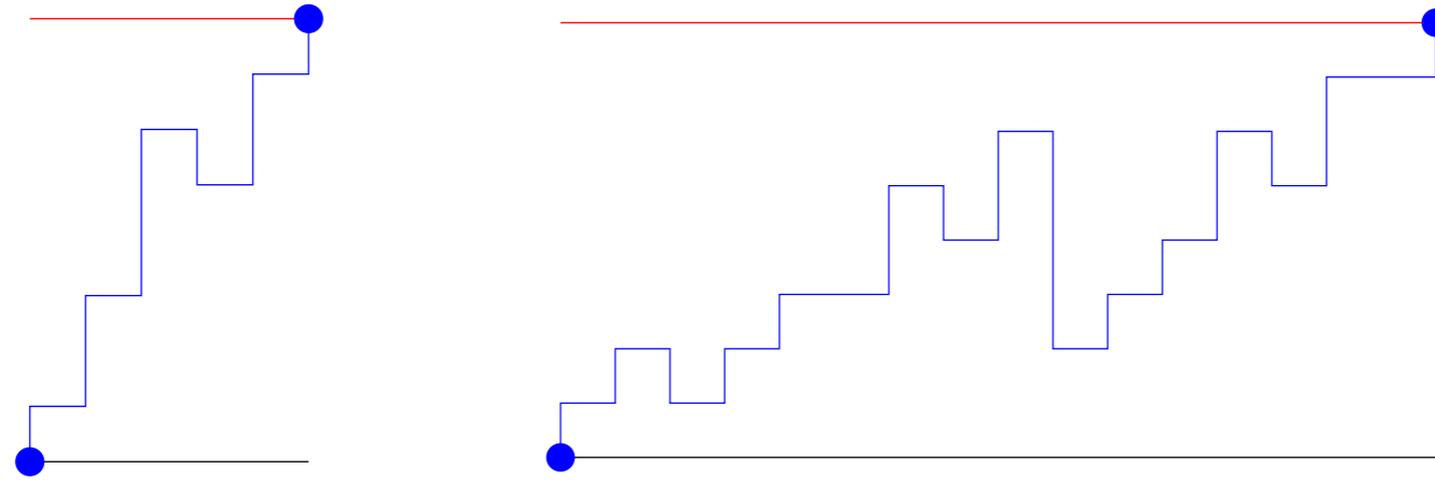
- Measure time between two successive events
- Heavily used in earthquake analysis
- Random distribution: both distribution of recurrence times, and cumulative distribution are exponential



$$p(t) = \tau^{-1} e^{-t/\tau} \quad P(t) = \int_t^{\infty} ds P(s) = e^{-t/\tau}$$

Very good agreement with random distribution!

First-Passage Processes



- Process by which a fluctuating quantity reaches a threshold for the first time
- **First-passage probability:** for the random variable to reach the threshold as a function of time.
- **Total probability:** that threshold is ever reached. May or may not equal 1
- **First-passage time:** the mean duration of the first-passage process. Can be finite or infinite

Marathon world record

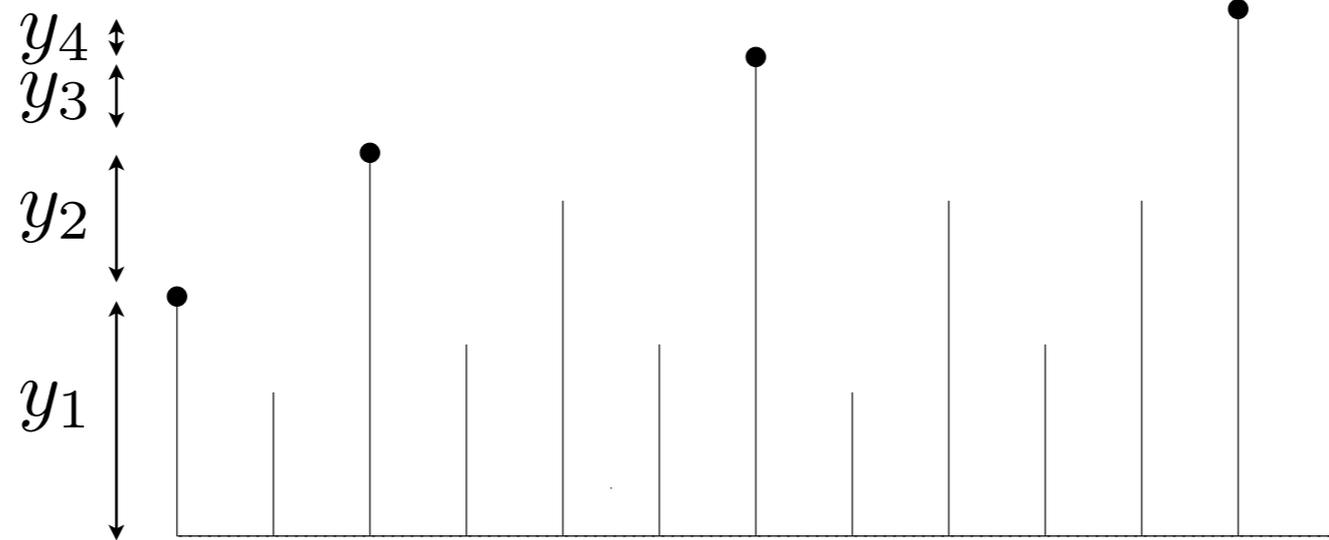
Year	Athlete	Country	Record	Improvement
2002	Khalid Khannuchi	USA	2:05:38	
2003	Paul Tergat	Kenya	2:04:55	0:43
2007	Haile Gebrselassie	Ethiopia	2:04:26	0:29
2008	Haile Gebrselassie	Ethiopia	2:03:59	0:27
2011	Patrick Mackau	Kenya	2:03:38	0:21
2013	Wilson Kipsang	Kenya	2:03:23	0:15

Incremental sequence of records

every record improves upon previous record by yet smaller amount

Are incremental sequences of records common?

Incremental Records



Incremental sequence of records

every record improves upon previous record by yet smaller amount

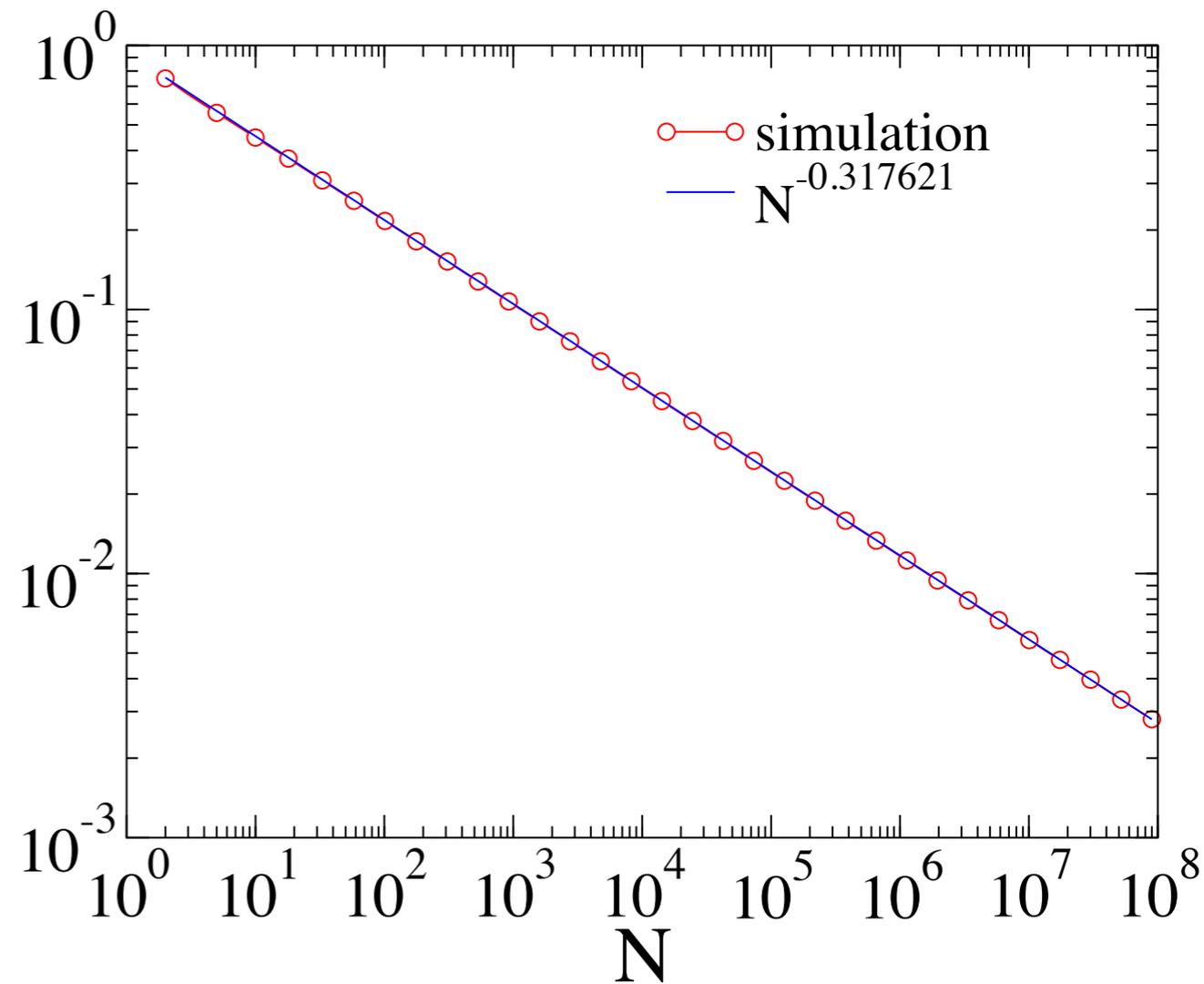
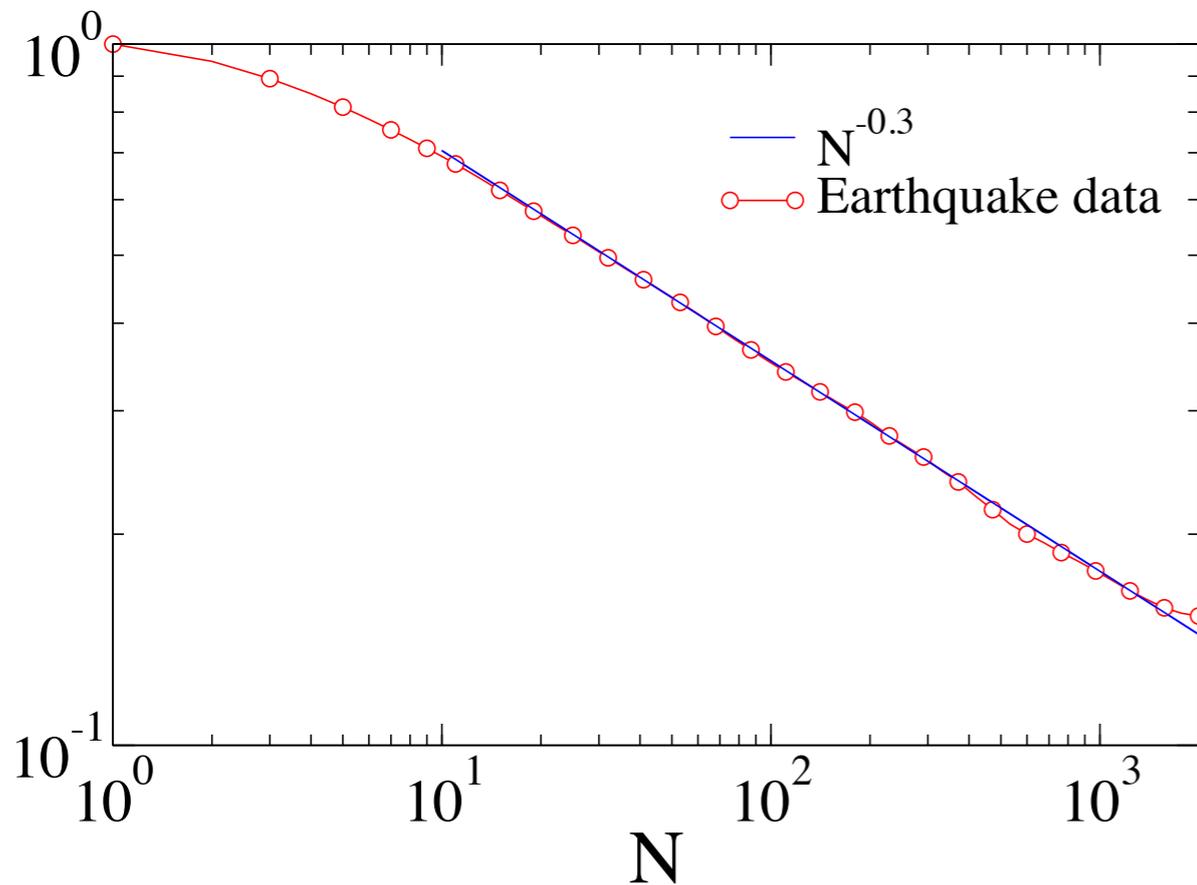
random variable = $\{0.4, 0.4, 0.6, 0.7, 0.5, 0.1\}$

latest record = $\{0.4, 0.4, 0.6, 0.7, 0.7, 0.7\}$ \uparrow

latest increment = $\{0.4, 0.4, 0.2, 0.1, 0.1, 0.1\}$ \downarrow

What is the probability all records are incremental?

Probability all records are incremental



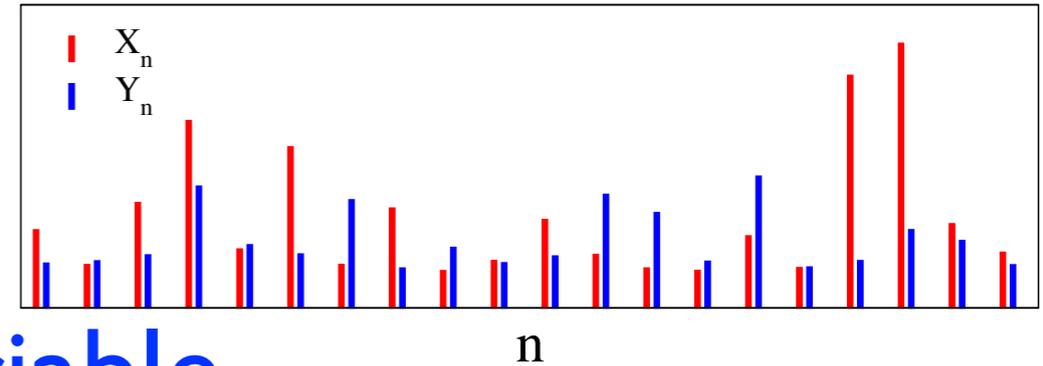
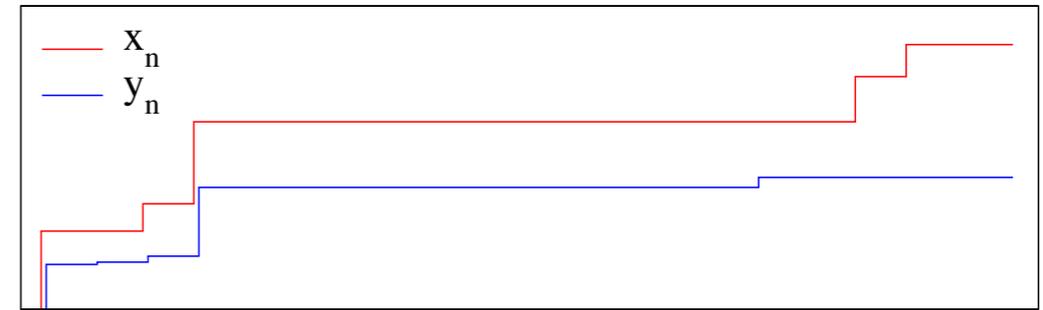
$$S_N \sim N^{-\nu} \quad \nu = 0.31762101$$

$$zg''(z) + (2 - \nu)g'(z) + e^{-z}g(z) = 0$$

Power law decay with nontrivial exponent
Problem is parameter-free

**II. Ordered records:
uncorrelated random variables**

Ordered Records



- Motivation: temperature records:
Record high increasing each year

- Two sequences of random variable

$$\{X_1, X_2, \dots, X_N\} \quad \text{and} \quad \{Y_1, Y_2, \dots, Y_N\}$$

- Independent and identically distributed variables

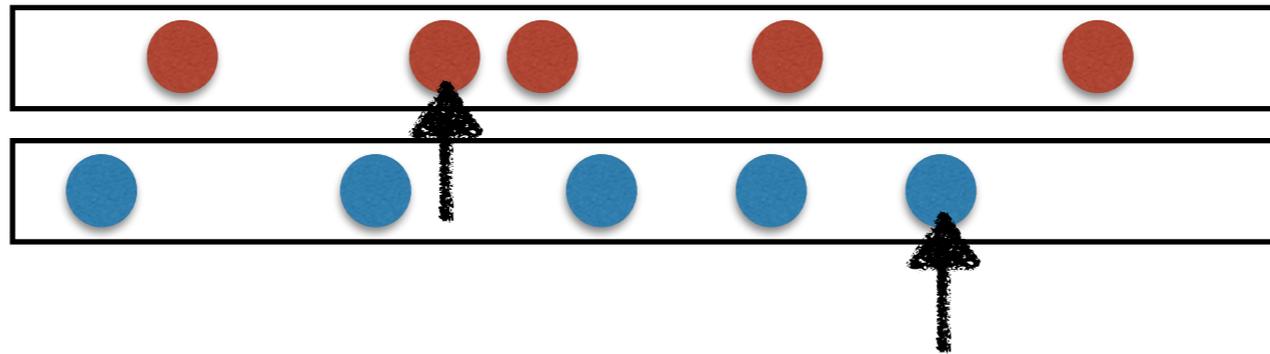
- Two corresponding sequences of records

$$x_n = \max\{X_1, X_2, \dots, X_n\} \quad \text{and} \quad y_n = \max\{Y_1, Y_2, \dots, Y_n\}$$

- Probability S_N records maintain perfect order

$$x_1 > y_1 \quad \text{and} \quad x_2 > y_2 \quad \cdots \quad \text{and} \quad x_N > y_N$$

Two Sequences



- Survival probability obeys closed recursion equation

$$S_N = S_{N-1} \left(1 - \frac{1}{2N} \right)$$

- Solution is immediate

$$S_N = \binom{2N}{N} 2^{-2N}$$

- Large-N: Power-law decay with rational exponent

$$S_N \simeq \pi^{-1/2} N^{-1/2}$$

Universal behavior: independent of parent distribution!

Ordered Random Variables

- Probability P_N variables are always ordered

$$X_1 > Y_1 \quad \text{and} \quad X_2 > Y_2 \quad \cdots \quad \text{and} \quad X_N > Y_N$$

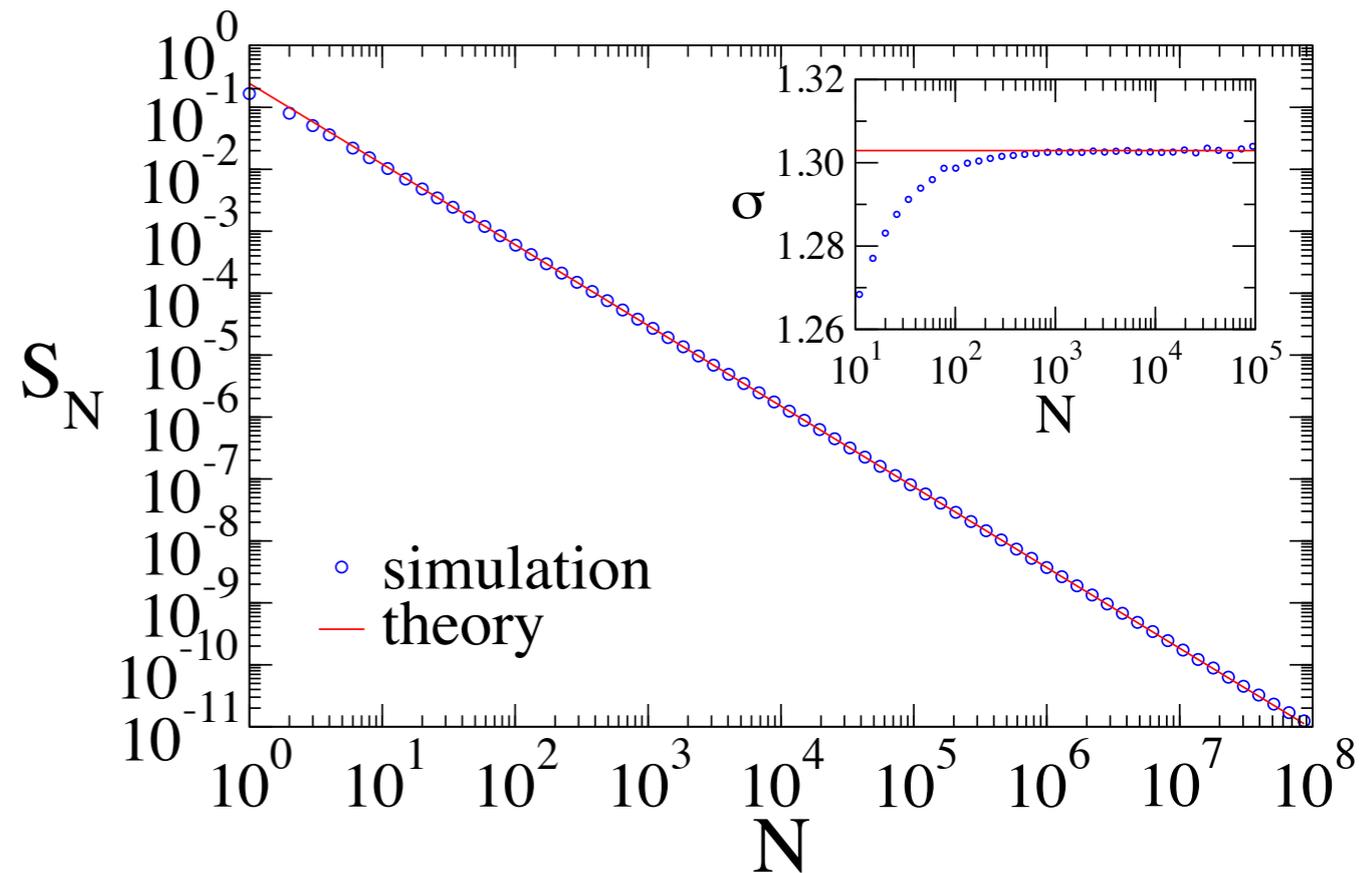
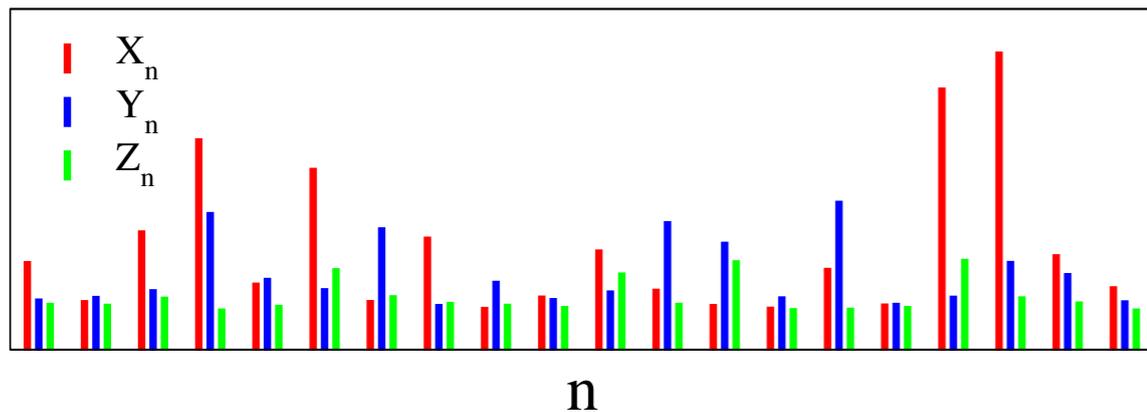
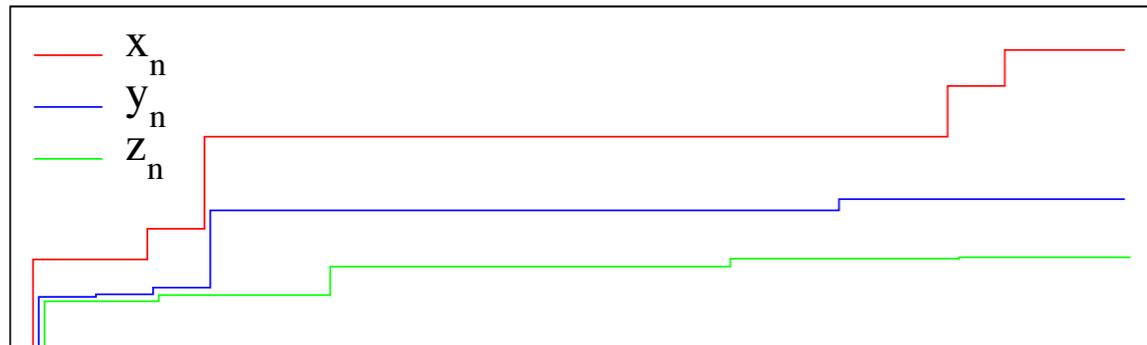
- Exponential decay

$$P_N = 2^{-N}$$

- Ordered records far more likely than ordered variables!
- Variables are uncorrelated
- Records are strongly correlated: each record “remembers” entire preceding sequence

Ordered records better suited for data analysis

Three sequences



- Third sequence of random variables

$$x_n > y_n > z_n \quad n = 1, 2, \dots, N$$

- Probability S_N records maintain perfect order
- Power-law decay with nontrivial exponent?

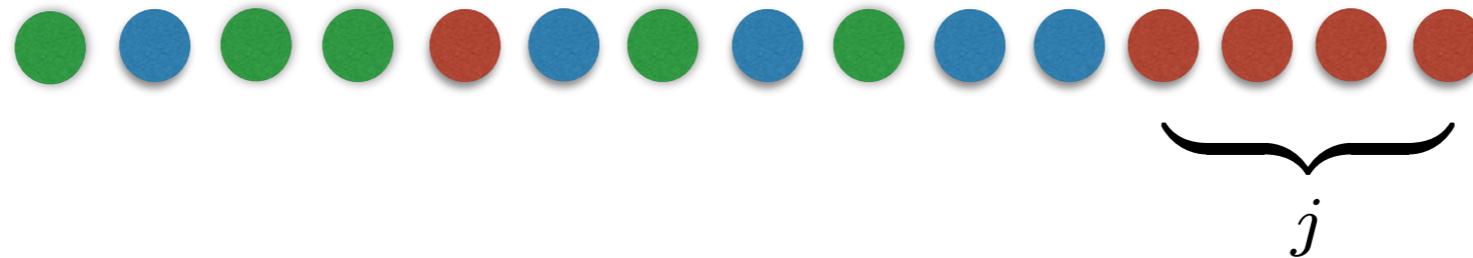
$$S_N \sim N^{-\sigma} \quad \text{with} \quad \sigma = 1.3029$$

Rank of median record

● leader

● median

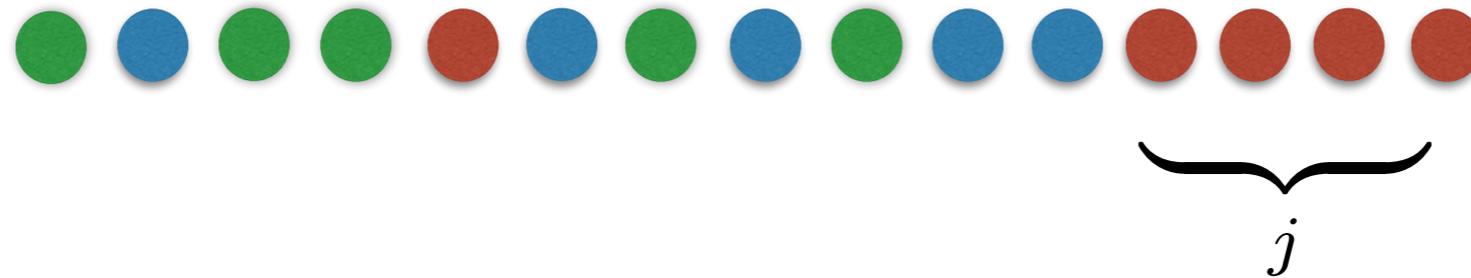
● laggard



- Closed equations for survival probability not feasible
- Focus on rank of the median record
- Rank of the trailing record irrelevant
- Joint probability $P_{N,j}$ that (i) records are ordered and (ii) rank of the median record equals j
- Joint probability gives the survival probability

$$S_N = \sum_{j=1}^N P_{N,j}$$

Closed Recursion Equations



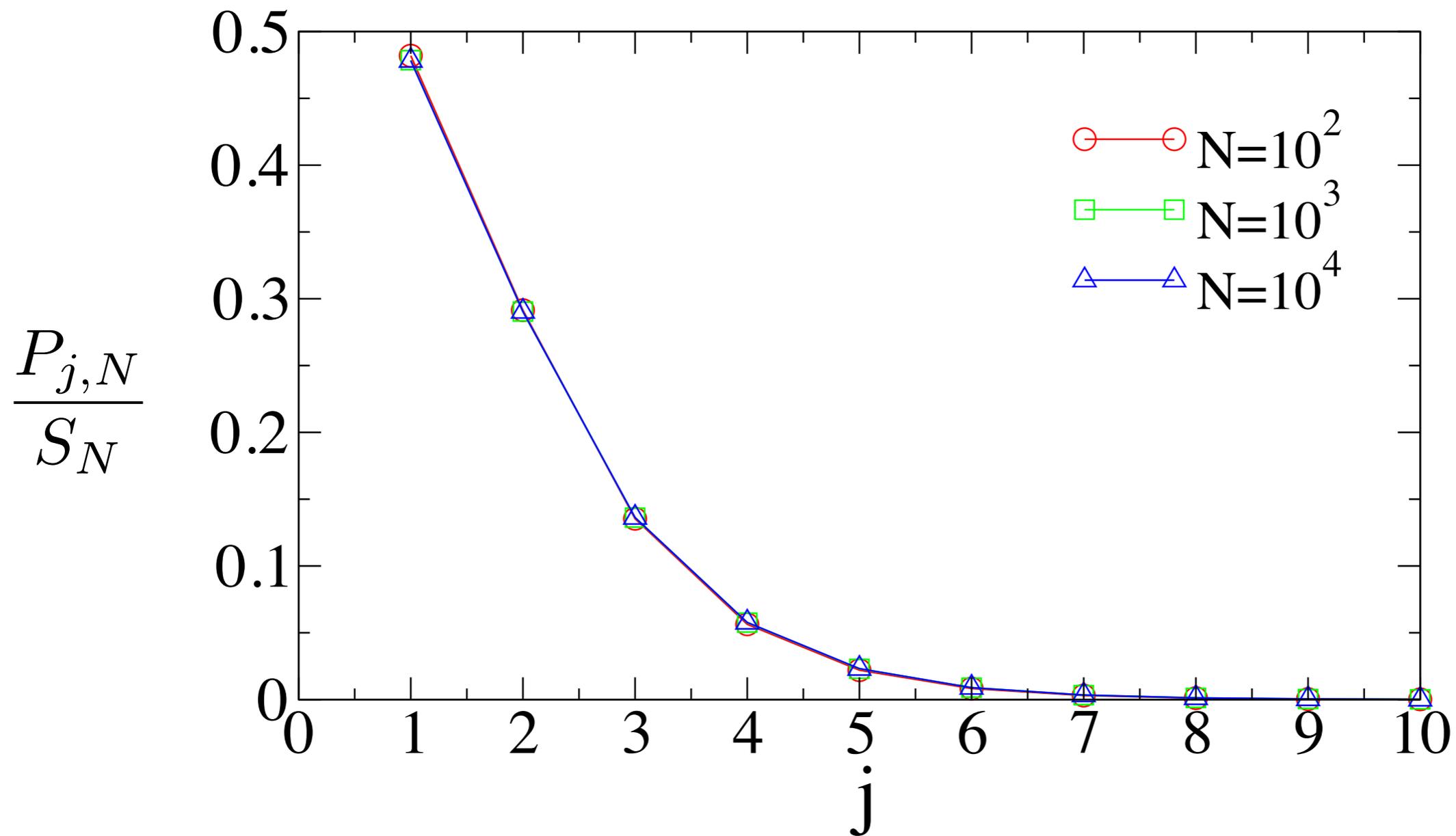
- Closed recursion equations for joint probability feasible

$$\begin{aligned}
 P_{N+1,j} = & \frac{3N+2-j}{3N+3} \frac{3N+1-j}{3N+2} \frac{3N-j}{3N+1} P_{N,j} \\
 & + \frac{3N+2-j}{3N+3} \frac{3N+1-j}{3N+2} \frac{j}{3N+1} P_{N,j-1} \\
 & + \frac{3N+2-j}{(3N+3)(3N+2)(3N+1)} \sum_{k=j}^{N+1} (3N-k) P_{N,k} \\
 & + \frac{3N+2-j}{(3N+3)(3N+2)(3N+1)} \sum_{k=j}^{N+1} k P_{N,k-1}
 \end{aligned}$$

- The survival probability for small N

N	S_N	$(3N)! S_N$
1	$\frac{1}{6}$	1
2	$\frac{29}{360}$	58
3	$\frac{4597}{90720}$	18388
4	$\frac{5393}{149688}$	17257600
5	$\frac{179828183}{6538371840}$	35965636600

Key Observation



Rank of median record j and N become uncorrelated!

Asymptotic Analysis

- Rank of median record j and N become uncorrelated!

$$P_{N,j} \simeq S_N p_j \quad \text{as } N \rightarrow \infty$$

- Assume power law decay for the survival probability

$$S_N \sim N^{-\sigma}$$

- The asymptotic rank distribution is normalized

$$\sum_{j=1}^{\infty} p_j = 1$$

- Rank distribution obeys a much simpler recursion

$$\sigma p_j = (j+1) p_j - \frac{j}{3} p_{j-1} - \frac{1}{3} \sum_{k=j}^{\infty} p_k$$

Scaling exponent σ is an eigenvalue

The Rank Distribution

- First-order differential equation for generating function

$$(3 - z) \frac{dP(z)}{dz} + P(z) \left(\frac{1}{1 - z} - \frac{3\sigma}{z} \right) = \frac{z}{1 - z} \quad P(z) = \sum_{j \geq 1} p_j z^{j+1}$$

- Solution

$$P(z) = \sqrt{\frac{1 - z}{3 - z}} \left(\frac{z}{3 - z} \right)^\sigma \int_0^z \frac{du}{(1 - u)^{3/2}} \frac{(3 - u)^{\sigma - 1/2}}{u^{\sigma - 1}}$$

- Behavior near $z=3$ gives tail of the distribution

$$p_j \sim j^{\sigma - 1/2} 3^{-j}$$

- Behavior near $z=1$ gives the scaling exponent

$${}_2F_1\left(-\frac{1}{2}, \frac{1}{2} - \sigma; \frac{3}{2} - \sigma; -\frac{1}{2}\right) = 0 \quad \implies \quad \sigma = 1.302931\dots$$

Three sequences: scaling exponent is nontrivial

Multiple Sequences

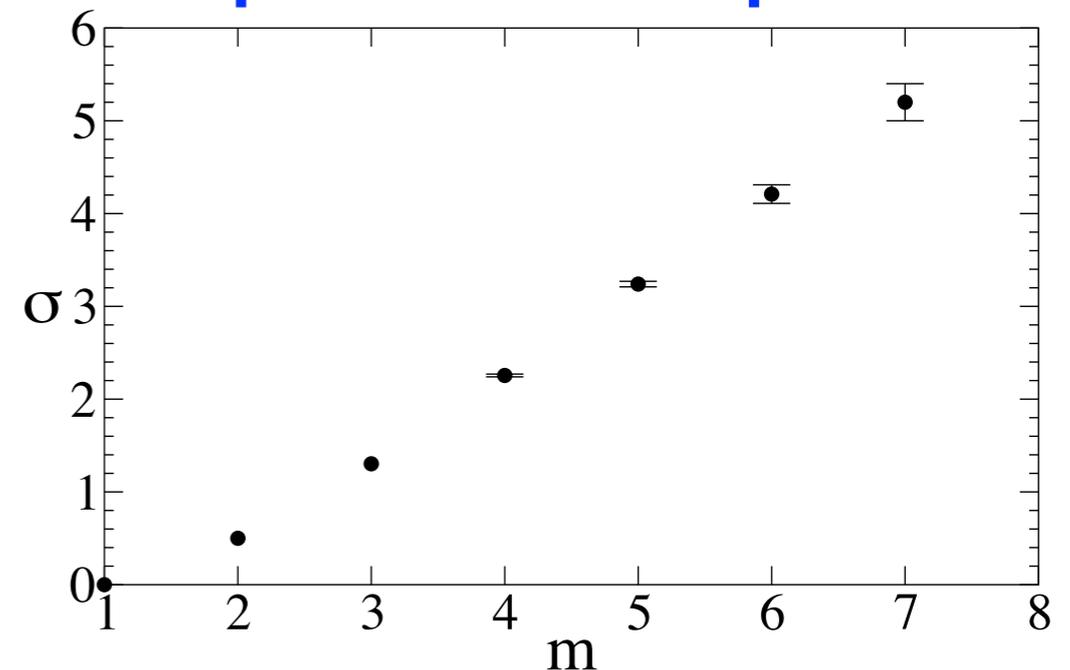
- Probability S_N that m records maintain perfect order
- Expect power-law decay with m -dependent exponent

$$S_N \sim N^{-\sigma_m}$$

- Lower and upper bounds

$$0 \leq m - \sigma_m \leq 1 + \frac{1}{2} + \frac{1}{3} + \dots + \frac{1}{m}$$

- Exponent grows linearly with number of sequences (up to a possible logarithmic correction)



$$\sigma_m \simeq m$$

In general, scaling exponent is nontrivial

Family of Ordering Exponents

- One sequence always in the lead: $1 \succ \text{rest}$

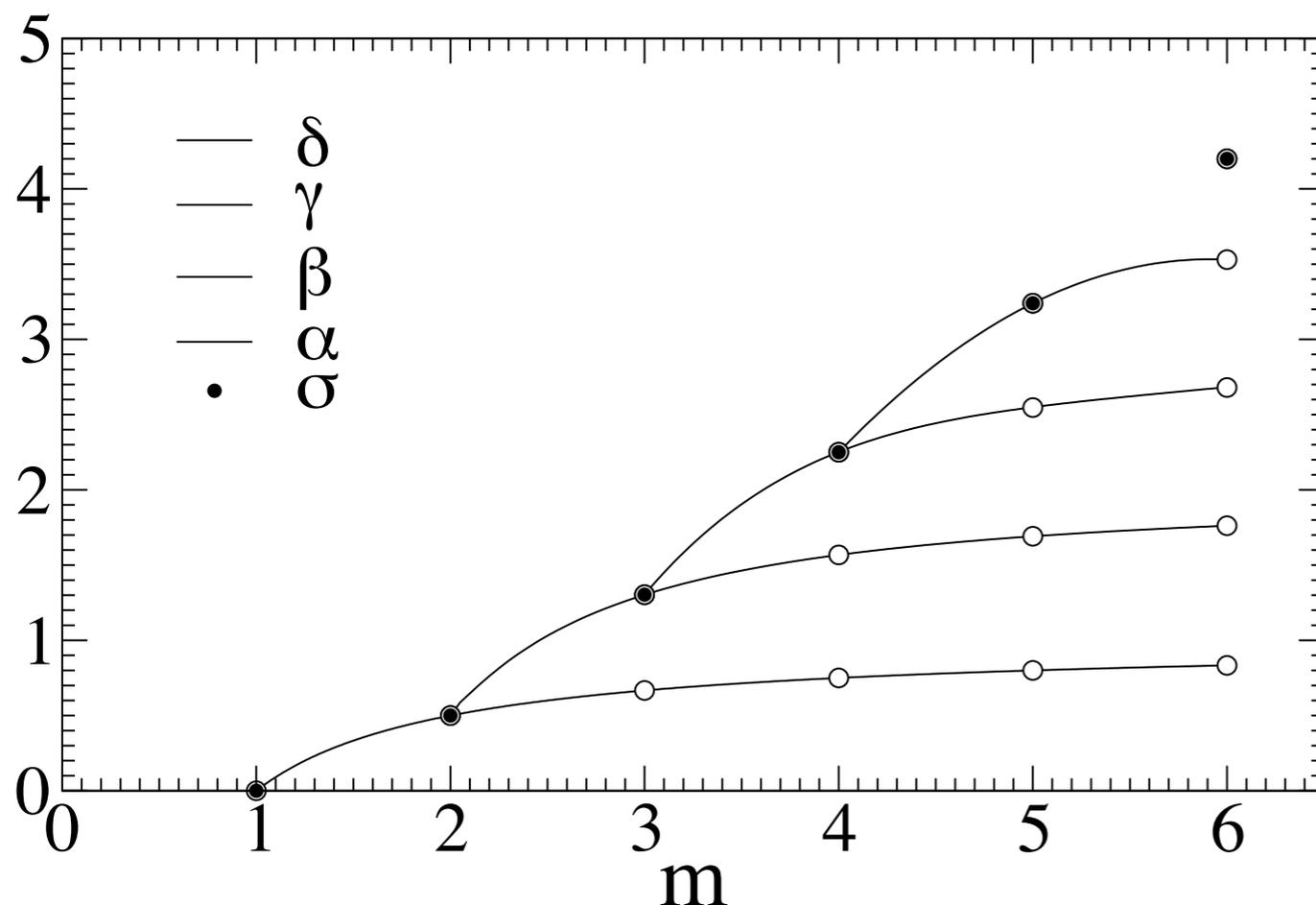
$$A_N \sim N^{-\alpha_m} \quad \alpha_m = 1 - \frac{1}{m}$$

- Two sequences always in the lead: $1 \succ 2 \succ \text{rest}$

$$B_N \sim N^{-\beta_m} \quad {}_2F_1\left(-\frac{1}{m-1}, \frac{m-2}{m-1} - \beta; \frac{2m-3}{m-1} - \beta; -\frac{1}{m-1}\right) = 0$$

- Three sequences always in the lead: $1 \succ 2 \succ 3 \succ \text{rest}$

$$C_N \sim N^{-\gamma_m}$$



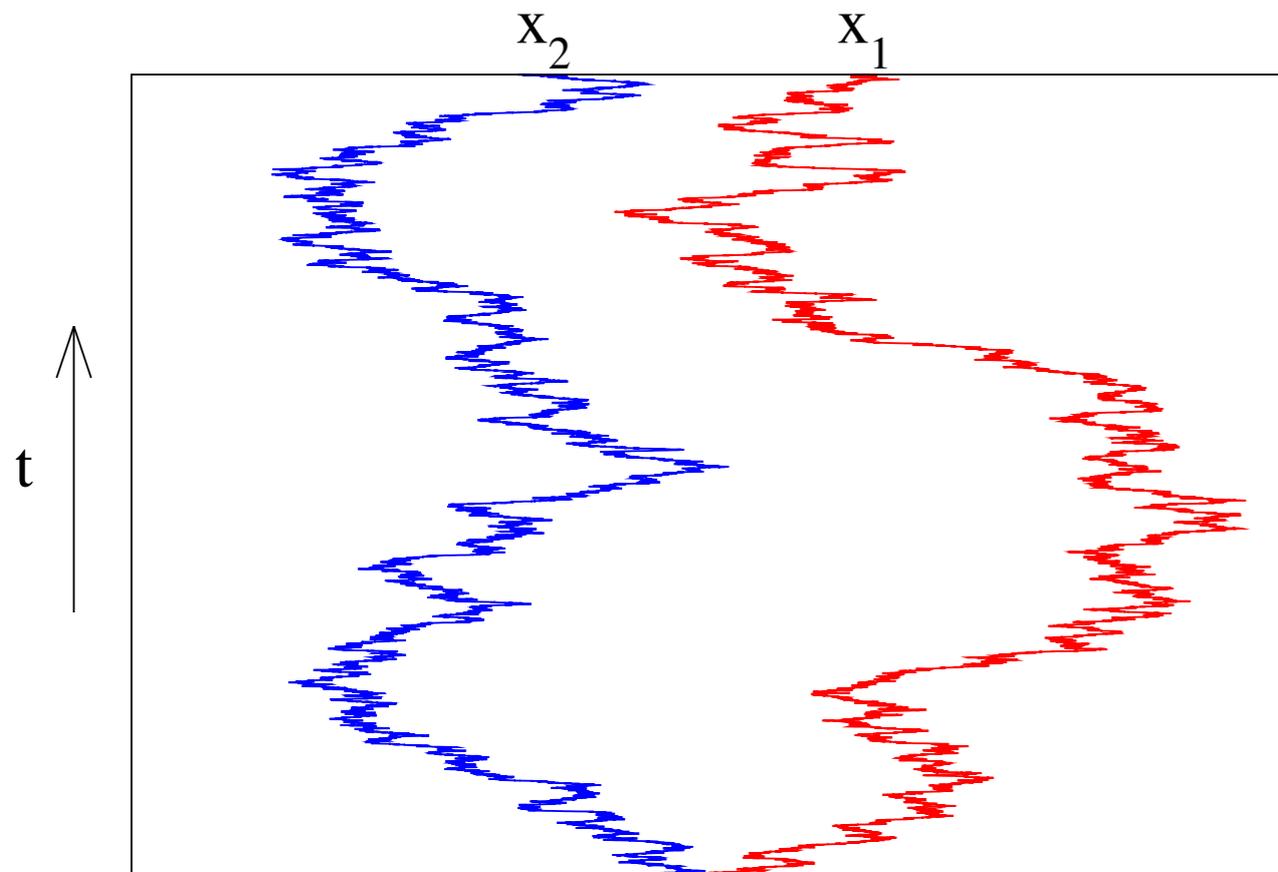
m	α	β	γ	δ
1	0			
2	1/2	1/2		
3	2/3	1.302931	1.302931	
4	3/4	1.56479	2.255	2.255
5	4/5	1.69144	2.547	3.24
6	5/6	1.76164	2.680	3.53

Summary I

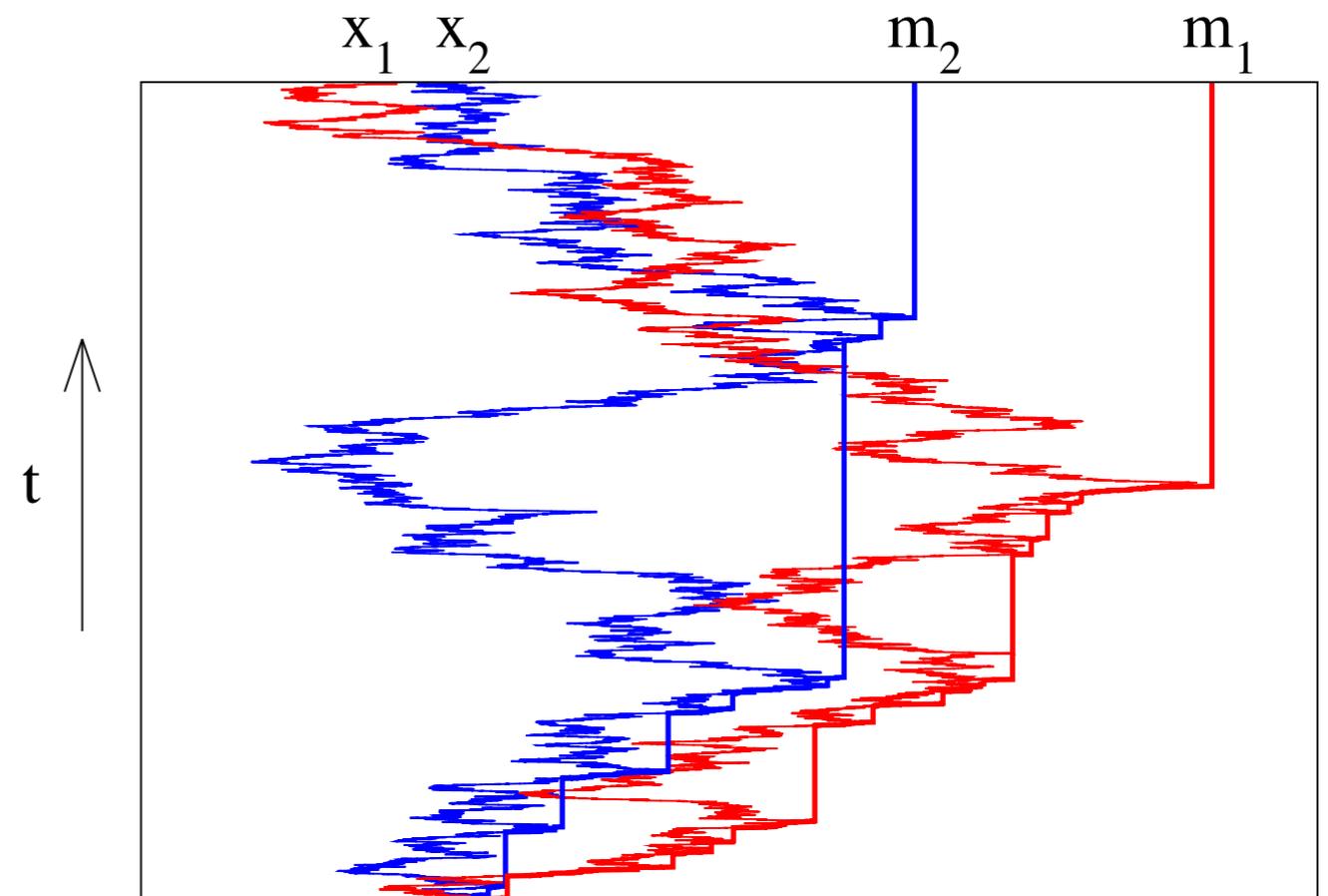
- Probability multiple sequences of records are ordered
- Uncorrelated random variables
- Survival probability independent of parent distribution
- Power-law decay with nontrivial exponent
- Exact solution for three sequences
- Scaling exponent grows linearly with number of sequences
- Key to solution: statistics of median record becomes independent of the sequence length (large N limit)
- Scaling methods allow us to tackle combinatorics

III. Ordered records: correlated random variables

Brownian Positions

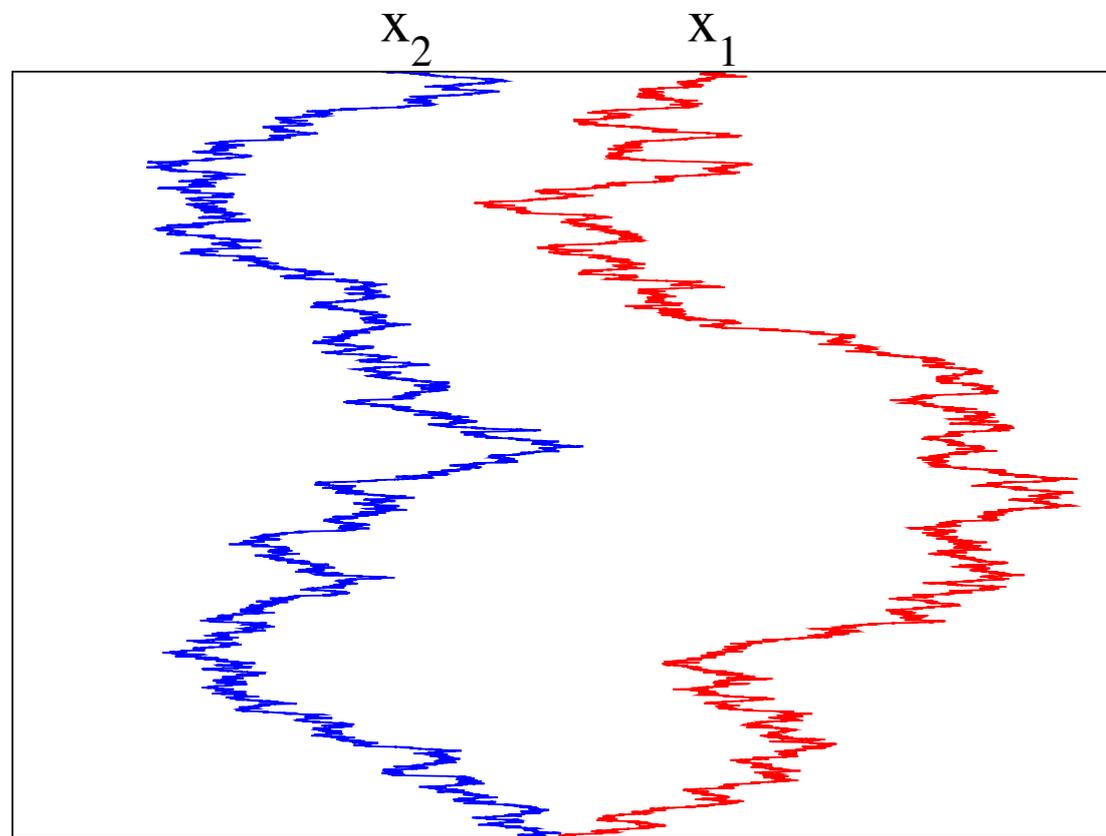


Brownian Records



First-Passage Kinetics: Brownian Positions

Probability two Brownian particle do not meet



- Universal probability Sparre Andersen 53

$$S_t = \binom{2t}{t} 2^{-t}$$

- Asymptotic behavior Feller 68

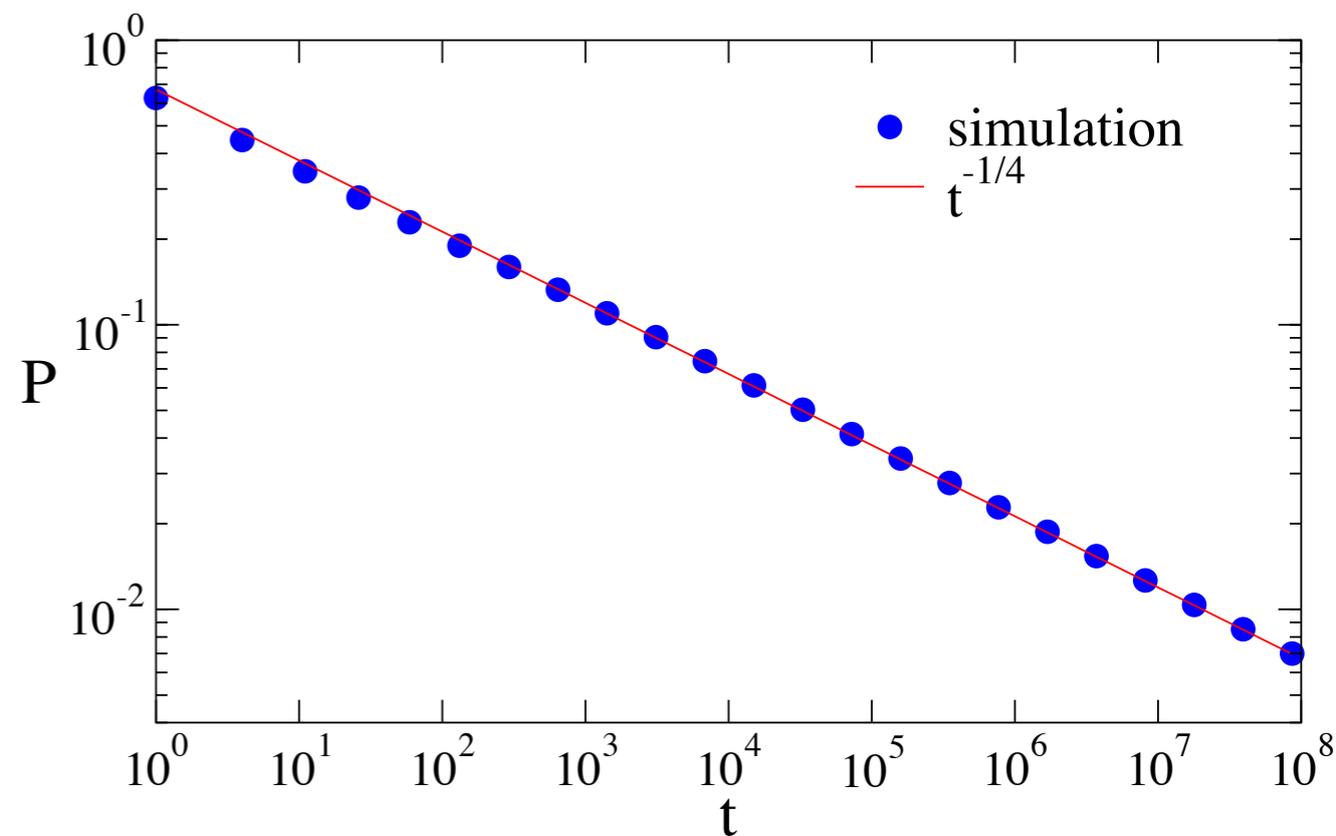
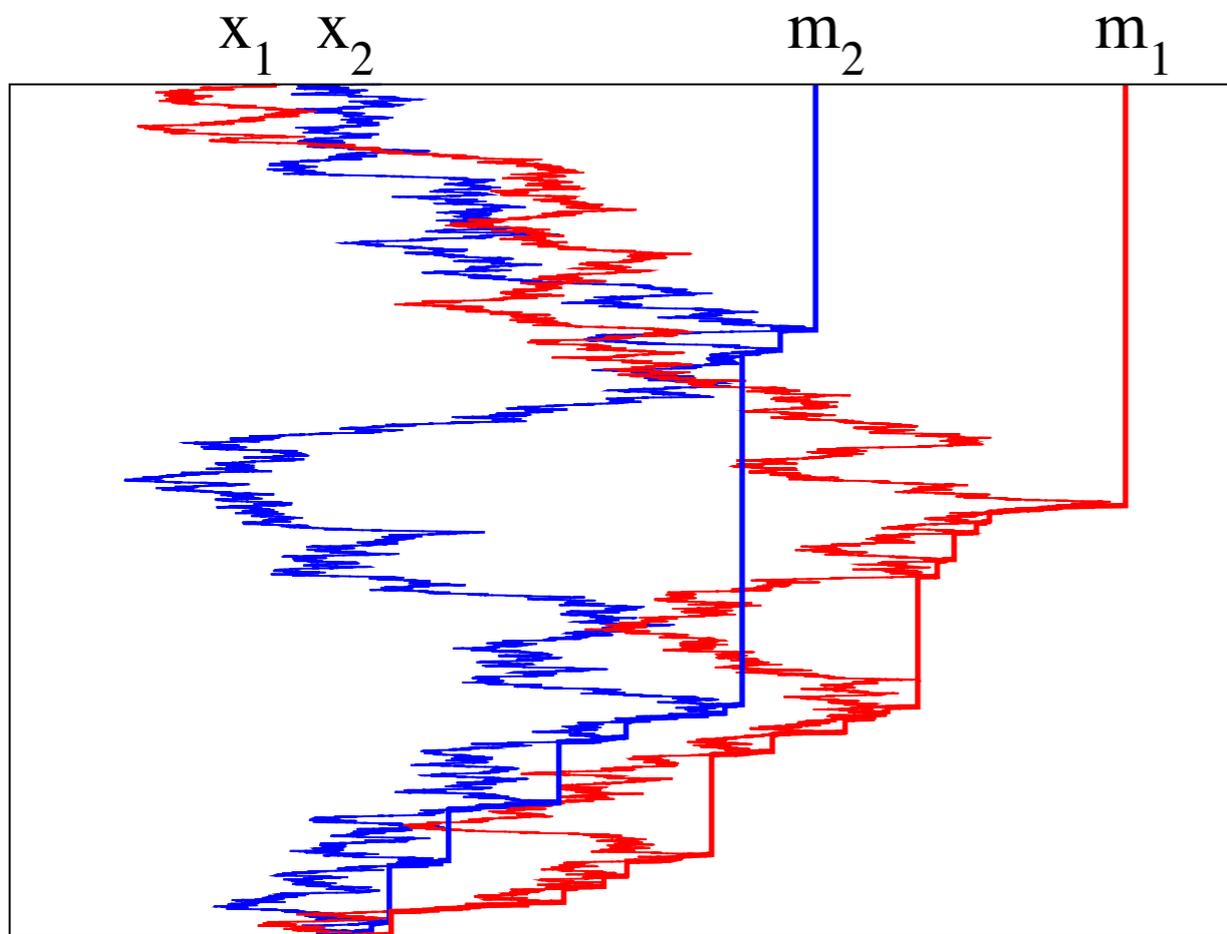
$$S \sim t^{-1/2}$$

Behavior holds for Levy flights, different mobilities, etc

Universal first-passage exponent

First-Passage Kinetics: Brownian Records

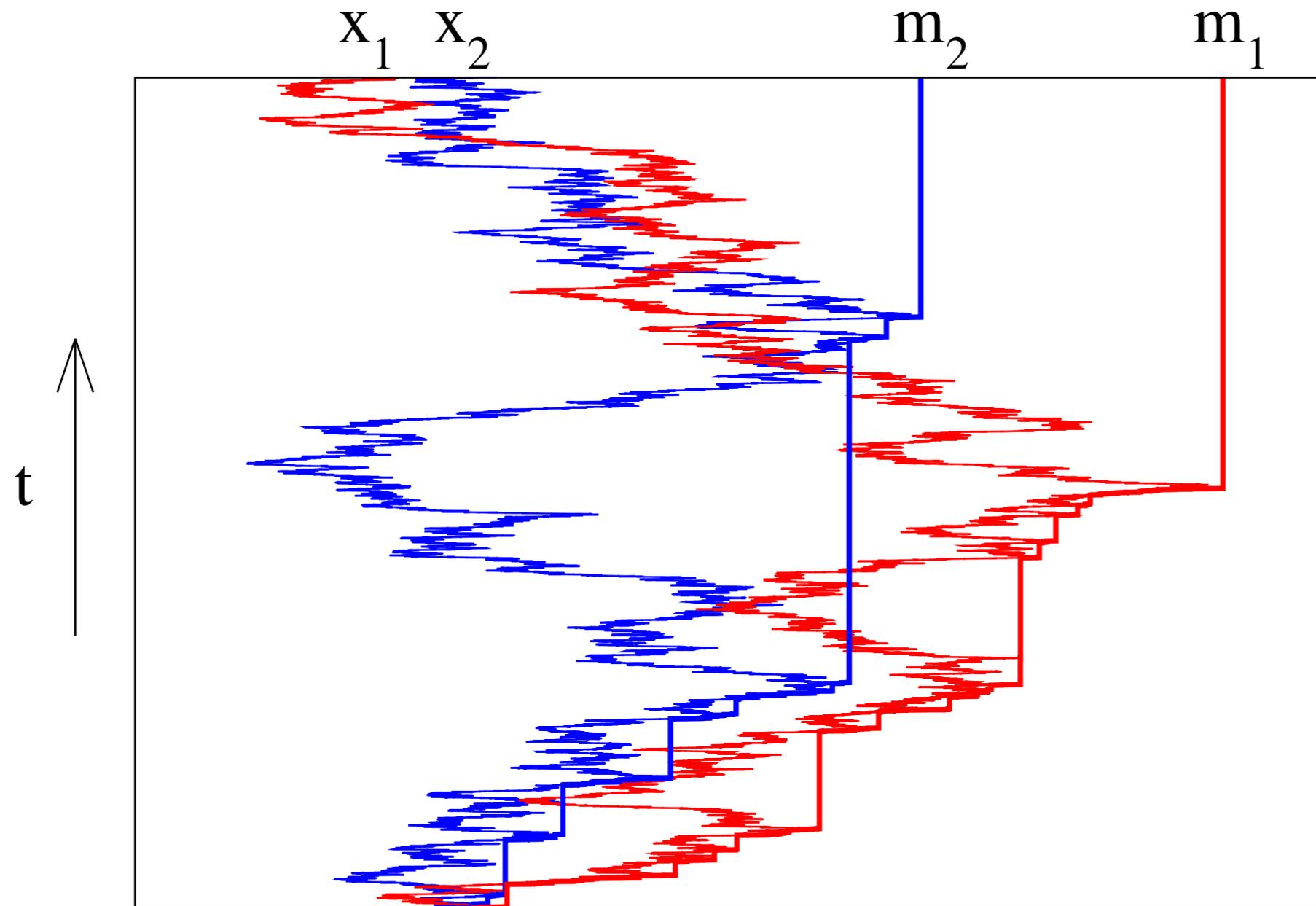
Probability running records remain ordered



$$S \sim t^{-\beta} \quad \beta = 0.2503 \pm 0.0005$$

Is $1/4$ exact? Is exponent universal?

$m_1 > m_2$ if and only if $m_1 > x_2$



From four variables to three

- Four variables: two positions, two records

$$m_1 > x_1 \quad \text{and} \quad m_2 > x_2$$

- The two records must always be ordered

$$m_1 > m_2$$

- Key observation: trailing record is irrelevant!

$$m_1 > m_2 \quad \text{if and only if} \quad m_1 > x_2$$

- Three variables: two positions, one record

$$m_1 > x_1 \quad \text{and} \quad m_1 > x_2$$

From three variables to two

- Introduce two distances from the record

$$u = m_1 - x_1 \quad \text{and} \quad v = m_1 - x_2$$

- Both distances undergo Brownian motion

$$\frac{\partial \rho(u, v, t)}{\partial t} = D \nabla^2 \rho(u, v, t)$$

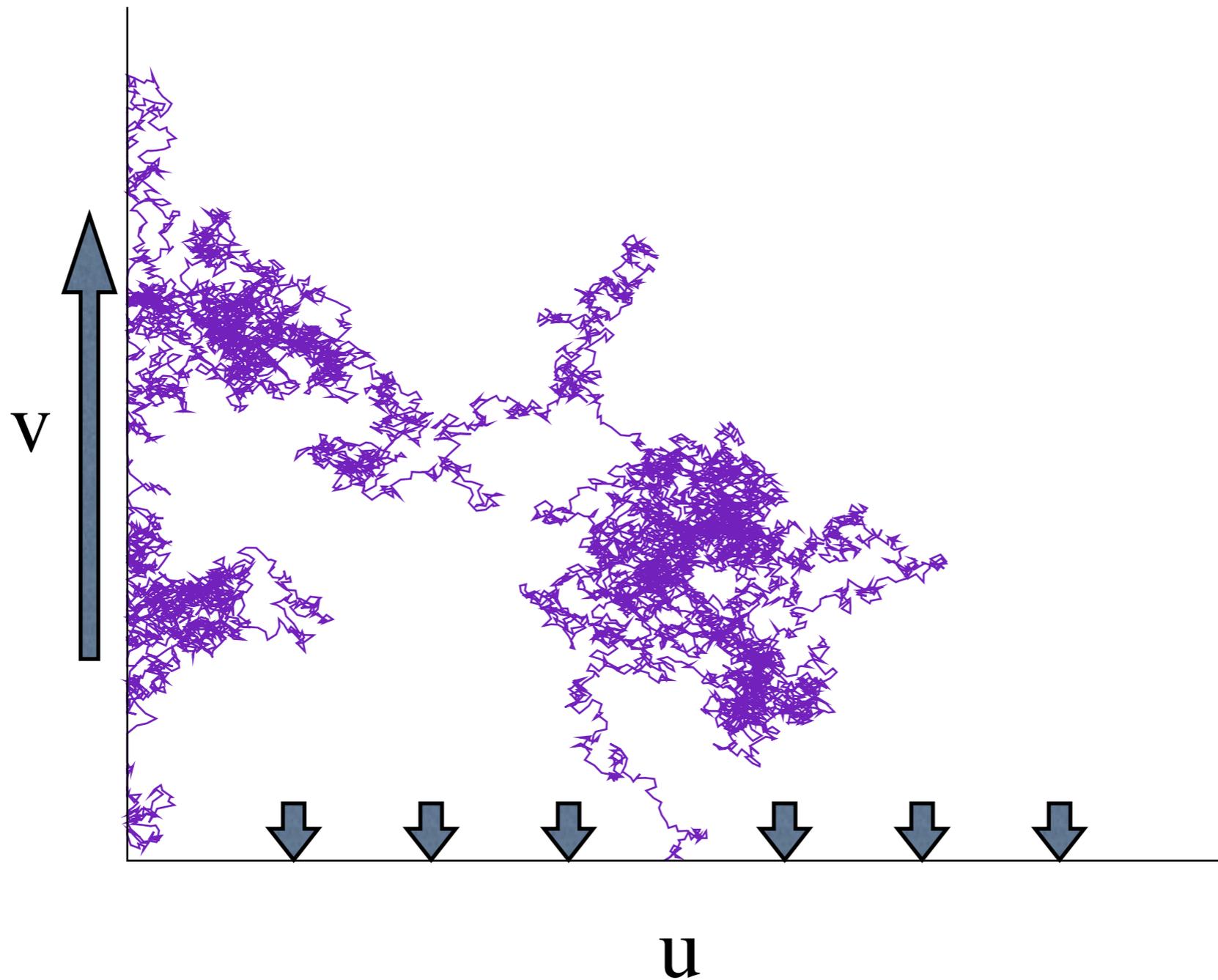
- Boundary conditions: (i) absorption (ii) advection

$$\rho|_{v=0} = 0 \quad \text{and} \quad \left(\frac{\partial \rho}{\partial u} - \frac{\partial \rho}{\partial v} \right) \Big|_{u=0} = 0$$

- Probability records remain ordered

$$P(t) = \int_0^\infty \int_0^\infty du dv \rho(u, v, t)$$

Diffusion in corner geometry



“Backward” evolution

- Study evolution as function of initial conditions

$$P \equiv P(u_0, v_0, t)$$

- Obeys diffusion equation

$$\frac{\partial P(u_0, v_0, t)}{\partial t} = D \nabla^2 P(u_0, v_0, t)$$

- Boundary conditions: (i) absorption (ii) advection

$$P|_{v_0=0} = 0 \quad \text{and} \quad \left(\frac{\partial P}{\partial u_0} + \frac{\partial P}{\partial v_0} \right) \Big|_{u_0=0} = 0$$

- Advection boundary condition is conjugate!

Solution

- Use polar coordinates

$$r = \sqrt{u_0^2 + v_0^2} \quad \text{and} \quad \theta = \arctan \frac{v_0}{u_0}$$

- Laplace operator

$$\nabla^2 = \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2}$$

- Boundary conditions: (i) absorption (ii) advection

$$P|_{\theta=0} = 0 \quad \text{and} \quad \left(r \frac{\partial P}{\partial r} - \frac{\partial P}{\partial \theta} \right) \Big|_{\theta=\pi/2} = 0$$

- Dimensional analysis + power law + separable form

$$P(r, \theta, t) \sim \left(\frac{r^2}{Dt} \right)^\beta f(\theta)$$

Selection of exponent

- Exponent related to eigenvalue of angular part of Laplacian

$$f''(\theta) + (2\beta)^2 f(\theta) = 0$$

- Absorbing boundary condition selects solution

$$f(\theta) = \sin(2\beta\theta)$$

- Advection boundary condition selects exponent

$$\tan(\beta\pi) = 1$$

- First-passage probability

$$P \sim t^{-1/4}$$

General Diffusivities

[ben Avraham](#)
[Leyvraz 88](#)

- Particles have diffusion constants D_1 and D_2

$$(x_1, x_2) \rightarrow (\hat{x}_1, \hat{x}_2) \quad \text{with} \quad (\hat{x}_1, \hat{x}_2) = \left(\frac{x_1}{\sqrt{D_1}}, \frac{x_2}{\sqrt{D_2}} \right)$$

- Condition on records involves ratio of mobilities

$$\sqrt{\frac{D_1}{D_2}} \hat{m}_1 > \hat{m}_2$$

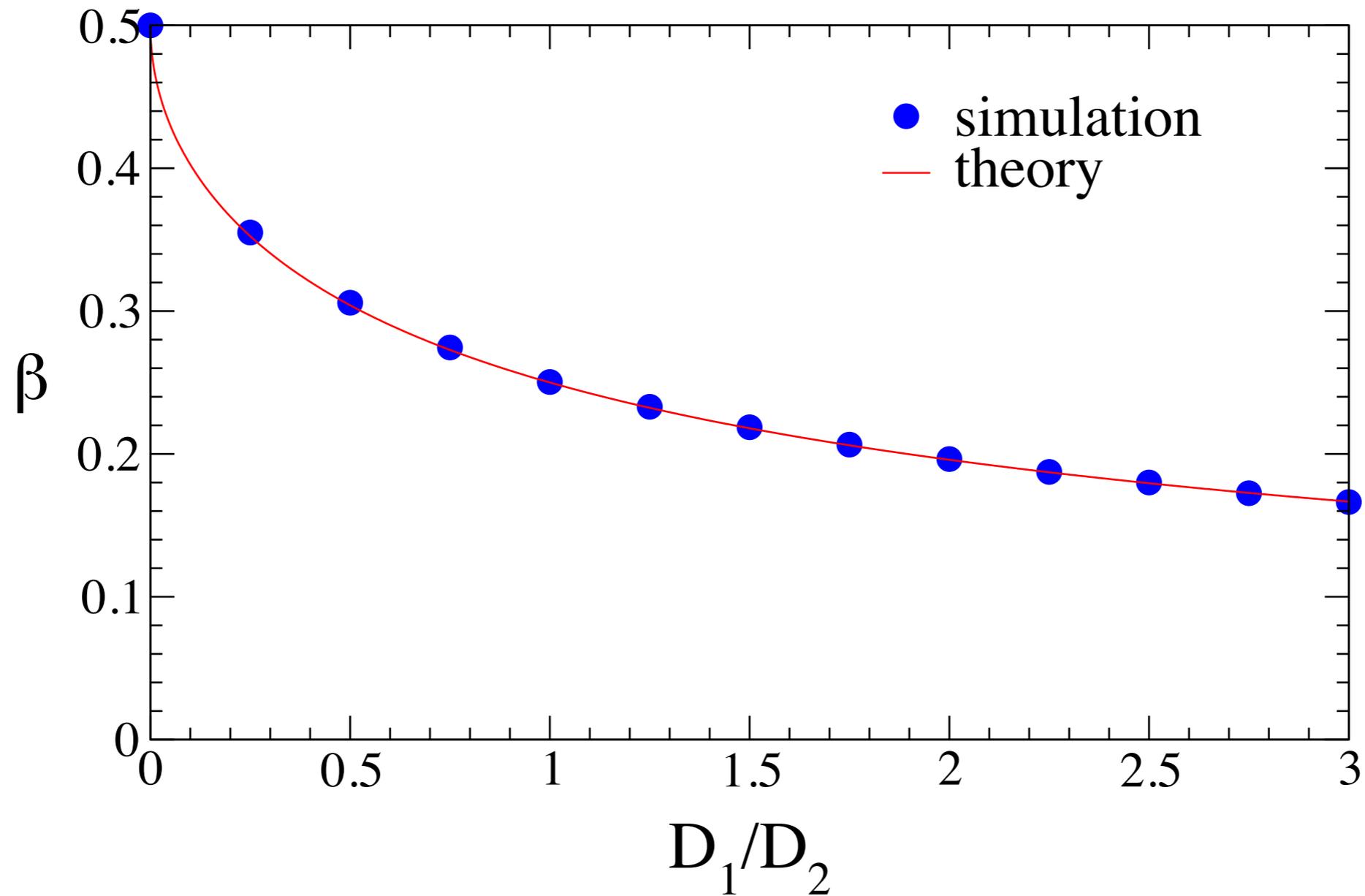
- Analysis straightforward to repeat

$$\sqrt{\frac{D_1}{D_2}} \tan(\beta\pi) = 1$$

- First-passage exponent: nonuniversal, mobility-dependent

$$\beta = \frac{1}{\pi} \arctan \sqrt{\frac{D_2}{D_1}}$$

Numerical verification



perfect agreement

Properties

- Depends on ratio of diffusion constants

$$\beta(D_1, D_2) \equiv \beta\left(\frac{D_1}{D_2}\right)$$

- Bounds: involve one immobile particle

$$\beta(0) = \frac{1}{2} \quad \beta(\infty) = 0$$

- Rational for special values of diffusion constants

$$\beta(1/3) = 1/3 \quad \beta(1) = 1/4 \quad \beta(3) = 1/6$$

- Duality: between “fast chasing slow” and “slow chasing fast”

$$\beta\left(\frac{D_1}{D_2}\right) + \beta\left(\frac{D_2}{D_1}\right) = \frac{1}{2}$$

Alternating kinetics: slow-fast-slow-fast

Multiple particles

- Probability n Brownian positions are perfectly ordered

$$P_n \sim t^{-\alpha_n} \quad \alpha_n = \frac{n(n-1)}{4}$$

Fisher & Huse 88

- Records perfectly ordered

$$m_1 > m_2 > m_3 > \dots > m_n$$

- In general, power-law decay

$$S_n \sim t^{-\nu_n}$$

n	ν_n	$\sigma_n/2$
2	1/4	1/4
3	0.653	0.651465
4	1.13	1.128
5	1.60	1.62
6	2.01	2.10

Uncorrelated variables provide an excellent approximation
Suggests some record statistics can be robust

Summary II

- First-passage kinetics of extremes in Brownian motion
- Problem reduces to diffusion in a two-dimensional corner with mixed boundary conditions
- First-passage exponent obtained analytically
- Exponent is continuously varying function of mobilities
- Relaxation is generally slower compared with positions
- Open questions: multiple particles, higher dimensions
- Why do uncorrelated variables represent an excellent approximation?

First-passage statistics of extreme values

- Survival probabilities decay as power law
- First-passage exponents are nontrivial
- Theoretical approach: differs from question to question
- Concepts of nonequilibrium statistical physics powerful: scaling, correlations, large system-size limit
- Many, many open questions
- Ordered records as a data analysis tool

Publications

1. Scaling Exponents for Ordered Maxima,
E. Ben-Naim and P.L. Krapivsky,
Phys. Rev. E **92**, 062139 (2015)
2. Slow Kinetics of Brownian Maxima,
E. Ben-Naim and P.L. Krapivsky,
Phys. Rev. Lett. **113**, 030604 (2014)
3. Persistence of Random Walk Records,
E. Ben-Naim and P.L. Krapivsky,
J. Phys. A **47**, 255002 (2014)
4. Scaling Exponent for Incremental Records,
P.W. Miller and E. Ben-Naim,
J. Stat. Mech. P10025 (2013)
5. Statistics of Superior Records,
E. Ben-Naim and P.L. Krapivsky,
Phys. Rev. E **88**, 022145 (2013)